



École des Ponts
ParisTech



UNSW
AUSTRALIA

Thermo-hydro-mechanical couplings and strain localization in Cosserat continuum

Application to stability analysis of rapid shear in faults

Hadrien Rattiez

Supervisors: Ioannis Stefanou, Jean Sulem

Acknowledgements: Manolis Veveakis, Thomas Poulet

UNIVERSITÉ
— PARIS — EST

Navier

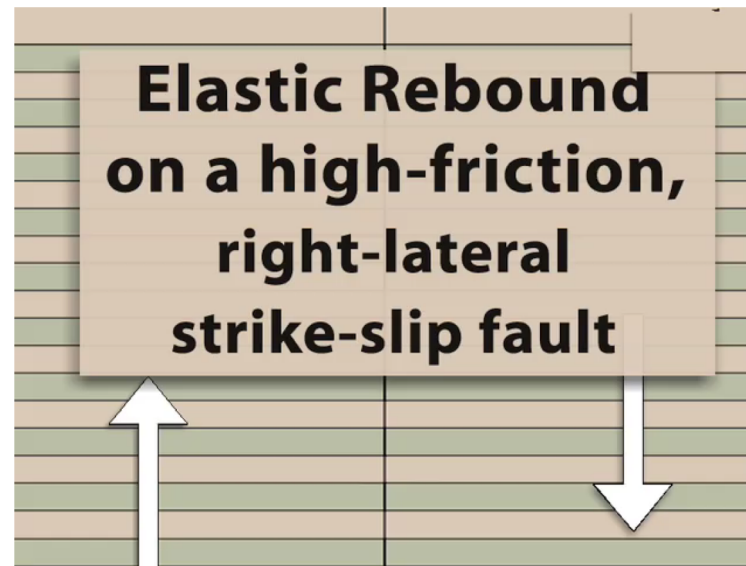


IFSTAR



Mechanism of earthquakes

Seismic faults



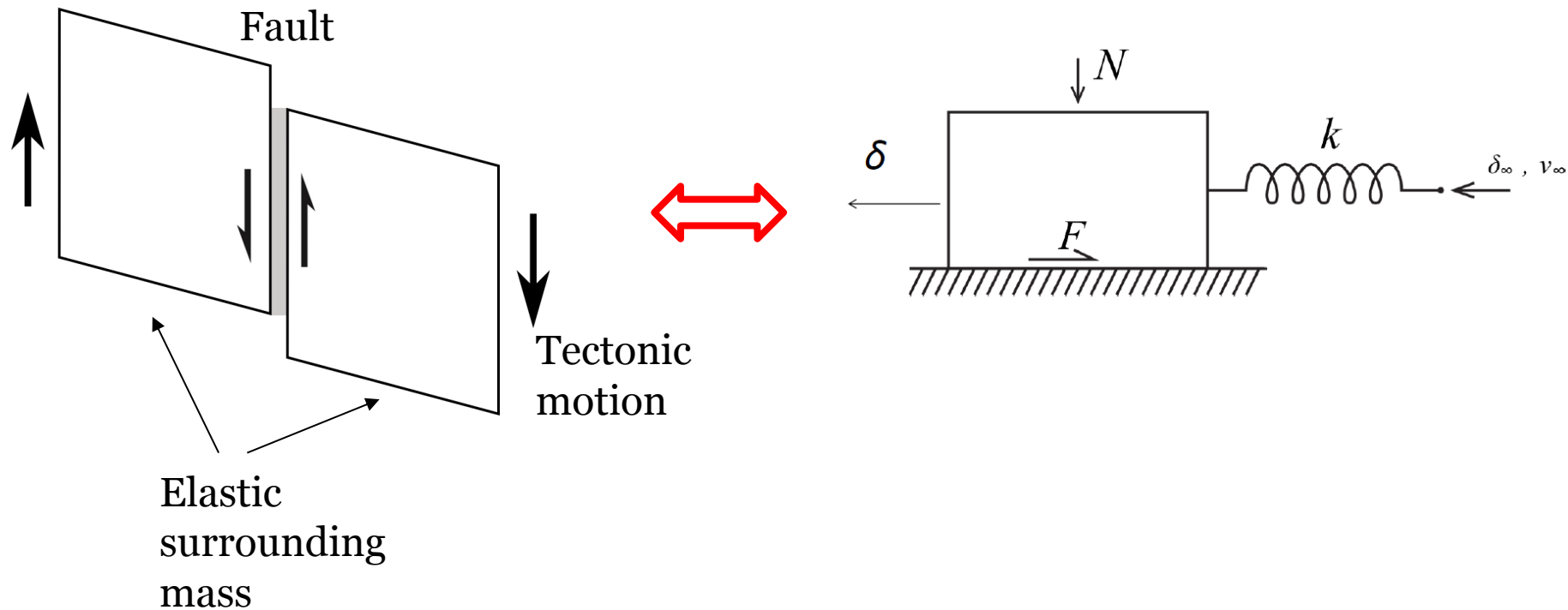
Animations from IRIS

“Earthquakes occur because fault strength weakens with increasing slip or slip rate. What physical processes determine how that weakening occurs?”

J. Rice, 2006

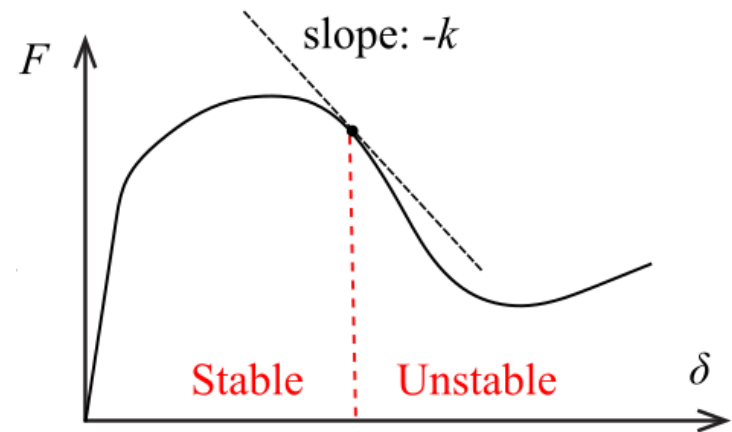
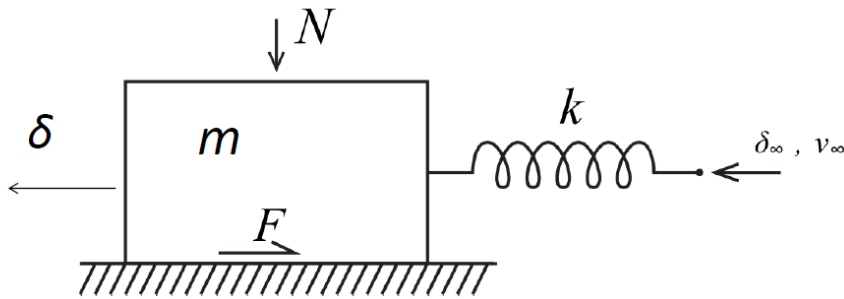
Spring-slider model

Frictional instabilities explained by a simple spring-slider model.



Stability of the spring-slider

- The friction law determines the stability of the system.



- Instability** condition:

$$\frac{dF}{d\delta} < -k$$



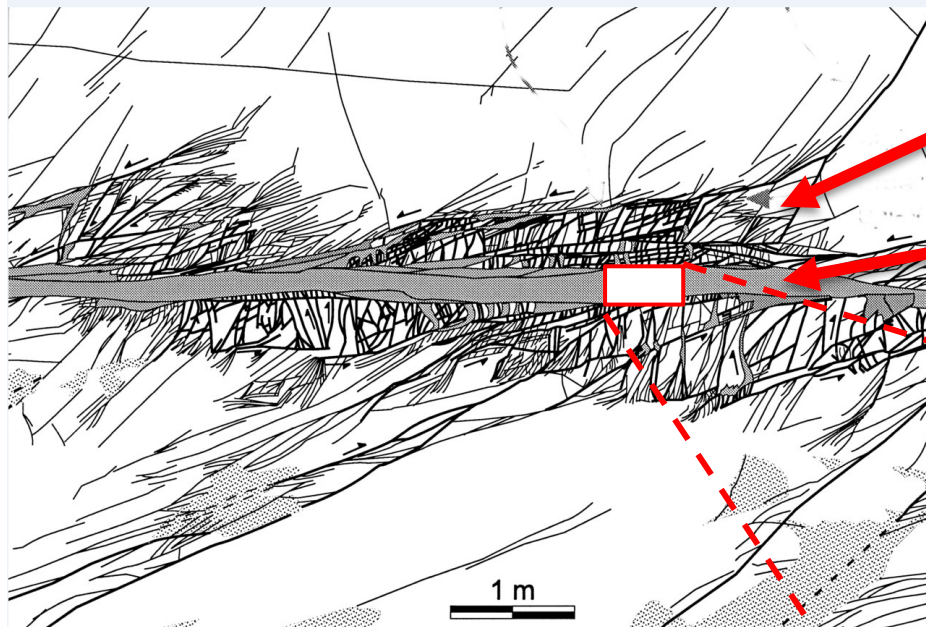
Importance of correctly capturing the softening behavior.

$$F(\delta, \dot{\delta}, \theta, \text{pore pressure, temperature, microstructure, ...})$$

Table of contents

- ❑ Fault mechanics and strain localization
- ❑ Mathematical modeling
- ❑ Linear Stability Analysis
- ❑ Finite Element simulations
- ❑ Conclusions and perspectives

Brittle fault zone



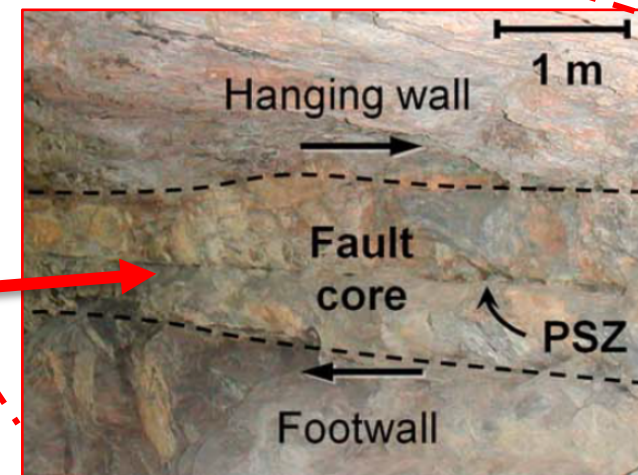
Myers et al. (1994)

Damaged zone
thickness: from ~10 m to ~1 km

Fault core
very fine crushed particles
thickness: from few μm to few mm

Principal Slip Zone (PSZ)

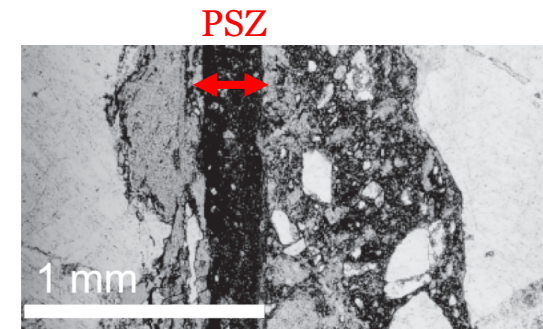
Poulet et al. (2014)



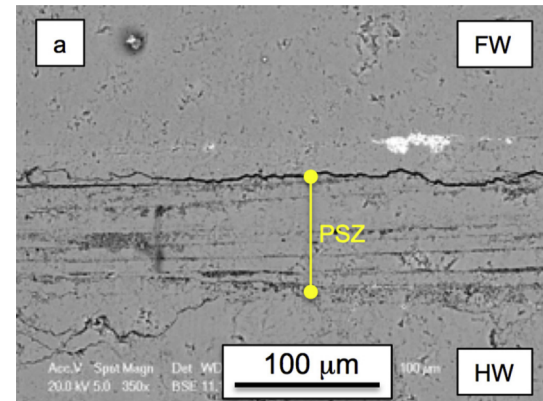
Field observation of Principal Slip Zones (PSZ)

- Sizes depend strongly on physical properties of the gouge.

Fault system	Thickness of the PSZ	Reference
Median Tectonic Line, Japan	3 mm	Wibberley et al., 2003
Chelungpu fault, China	50-300 μm	Heermance et al., 2003
Longmenshan fault, China	1cm	Li et al. , 2013
Punchbowl fault, USA	100-300 μm	Chester et al., 2003
Northern Apennines, Italy	10-40 μm	De Paola et al., 2008



PSZ in Nevada
(Shipton et al., 2006)



PSZ in M. Maggio, Italy
(Collettini et al., 2014)

Weakening and localization

Major role of the width of the slip zone:

- In the energy budget of the system: control of the frictional heating.
- In the stability of the fault (stronger weakening for thinner shear zones).

Evolution of the width of the slip zone:

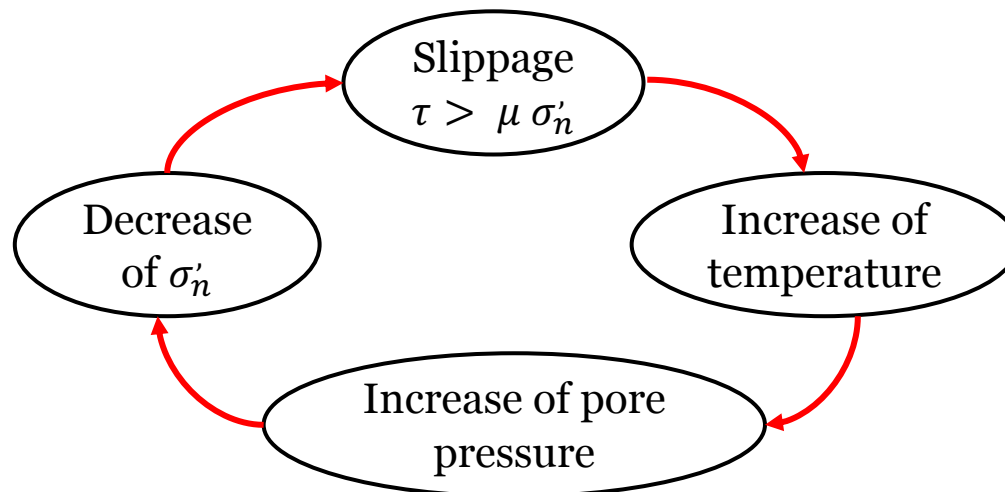
Stronger weakening favors a decrease of the localized zone thickness, heat and fluid diffusion tend to broaden it.

Constitutive models that can describe finite width of the slip zone:

- Rate dependent constitutive laws ([Rice, 2006](#), [Platt et al., 2014](#)).
- Higher order continuum theories ([Sulem et al., 2011](#), [Sulem & Stefanou, 2016](#)).

Examples of weakening and multiphysical couplings

- **Thermal decomposition of minerals** : dehydration of clay minerals (*Brantut et al., 2008*), decomposition of carbonates (*Sulem & Famin, 2009, Collettini et al., 2014*).
- **Flash heating** and shear weakening at micro-asperity contacts (*Rice, 1999, 2006, Spagnuolo et al., 2016, Brantut & Viesca 2017*).
- **Thermal pressurization of pore fluids** (*Lachenbruch, 1980, Sulem et al. 2005, Rice 2006, Viesca & Garagash 2015*)

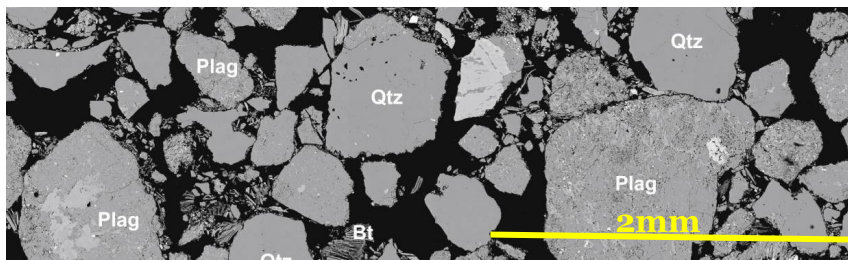


Importance of the grain size

- **Grain size reduction** inside strain localization zones (*Karato, 2008*)
- Effects of grain size on **physical processes** (poromechanics, friction, chemical interactions, etc...)

⇒ Need for a theory that takes into account the **size** of the microstructure and its evolution.

Host rock



Shear band

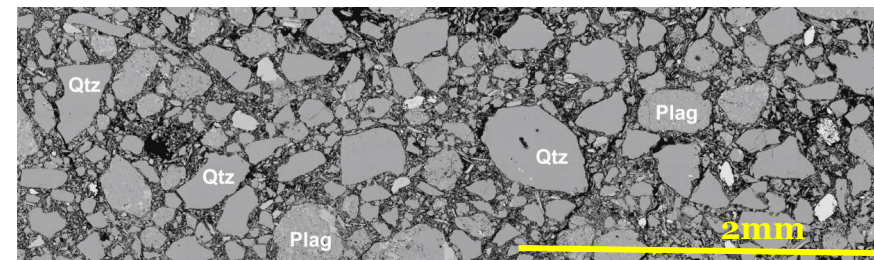
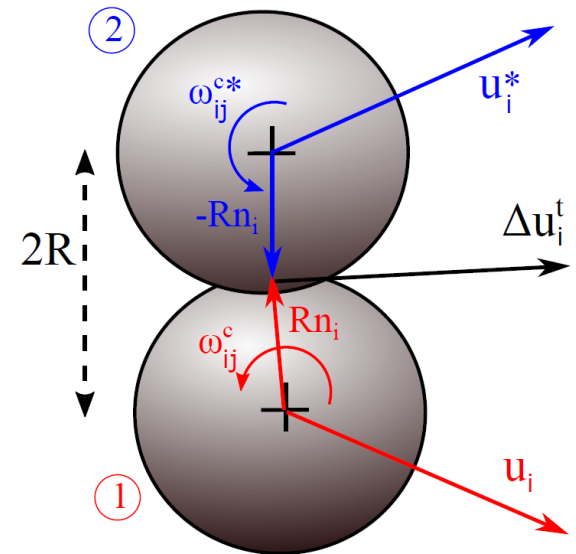
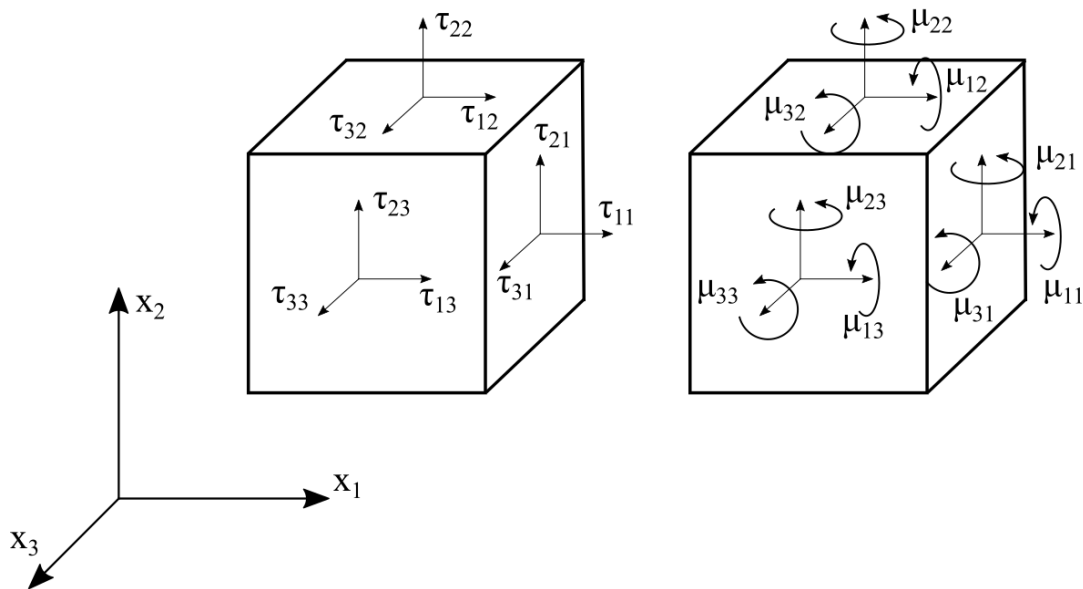


Table of contents

- ❑ Fault mechanics and strain localization
- ❑ **Mathematical modeling**
- ❑ Linear Stability Analysis
- ❑ Finite Element simulations
- ❑ Conclusions and perspectives

Cosserat continuum theory

- **Granular** character of fault gouges.
- Additional dofs: **rotations**.
- **Internal lengths** related to the grain size.



Cosserat continuum plasticity

- Drucker-Prager yield surface ([Mühlhaus & Vardoulakis, 1987](#)) with hardening:

$$F = \tau + \mu\sigma' - c$$

$$H_s = -\frac{d\mu}{d\gamma^p}\sigma'$$

- Generalized stress and strain invariants ([Sulem and Vardoulakis, 1990](#)):

$$\tau = \sqrt{h_1 s_{ij} s_{ij} + h_2 s_{ij} s_{ji} + \frac{1}{R^2} (h_3 \mu_{ij} \mu_{ij} + h_4 \mu_{ij} \mu_{ji})}$$

$$\gamma^p = \sqrt{g_1 e_{ij}^p e_{ij}^p + g_2 e_{ij}^p e_{ji}^p + R^2 (g_3 \kappa_{ij}^p \kappa_{ij}^p + g_4 \kappa_{ij}^p \kappa_{ji}^p)}$$

s_{ij} and e_{ij}^p are, respectively, the deviatoric part of the stress and the plastic strain.

R is the internal length.

h_i and g_i are coefficients.

Governing equations

A fault zone is modelled as an infinite layer under shear.

Momentum balance:

$$\tau_{ij,j} - \rho \frac{\partial^2 U_i}{\partial t^2} = 0$$

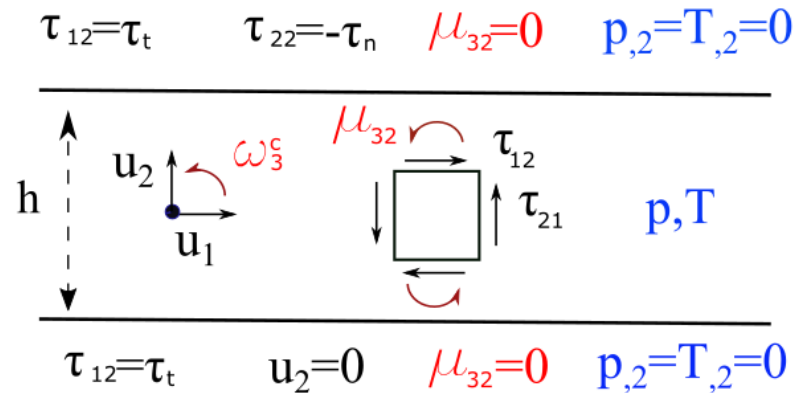
$$\mu_{ij,j} - e_{ijk} \tau_{jk} - \rho I \frac{\partial^2 \omega_i^c}{\partial t^2} = 0$$

Elasto-plastic constitutive law:

$$\dot{\tau}'_{ij} = C_{ijkl}^{ep} \dot{\gamma}_{kl} + D_{ijkl}^{ep} \dot{\kappa}_{kl} + E_{ijkl}^{ep} \dot{T} \delta_{kl}$$

$$\dot{\mu}_{ij} = M_{ijkl}^{ep} \dot{\kappa}_{kl} + L_{ijkl}^{ep} \dot{\gamma}_{kl} + N_{ijkl}^{ep} \dot{T} \delta_{kl}$$

Terzaghi effective stress: $\tau'_{ij} = \tau_{ij} + p \delta_{ij}$



Energy balance:

$$\frac{\partial T}{\partial t} - c_{th} \frac{\partial^2 T}{\partial z^2} = \frac{1}{\rho C} \underbrace{(\sigma_{ij} \dot{\epsilon}_{ij}^p + \mu_{ij} \dot{\kappa}_{ij}^p)}_{\text{Plastic work}}$$

Mass balance:

$$\frac{\partial p}{\partial t} = c_{hy} p_{,ii} + \underbrace{\frac{\lambda^*}{\beta^*} \frac{\partial T}{\partial t}}_{\text{Thermal pressurisation}} - \underbrace{\frac{1}{\beta^*} \frac{\partial \epsilon_v}{\partial t}}_{\text{Porosity variation}}$$

Table of contents

- ❑ Fault mechanics and strain localization
- ❑ Mathematical modeling
- ❑ **Linear Stability Analysis**
- ❑ Finite Element simulations
- ❑ Conclusions and perspectives

Linear stability analysis

- Perturbation of the steady state of homogeneous deformation, temperature and pore pressure:

$$\bar{T}, \bar{\gamma}_{ij}, \bar{p}$$



$$T(z, t) = \bar{T} + T^*(z, t)$$

$$\gamma_{ij}(z, t) = \bar{\gamma}_{ij} + \gamma_{ij}^*(z, t)$$

$$p(z, t) = \bar{p} + p^*(z, t)$$

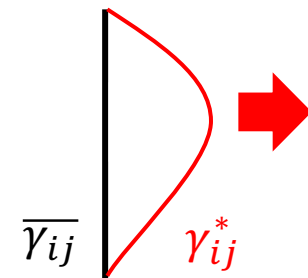
$$T^*(z, t) = \Theta \exp(s t) \exp(2\pi i \frac{z}{\lambda})$$

$$\gamma_{ij}^*(z, t) = E_{ij} \exp(s t) \exp(2\pi i \frac{z}{\lambda})$$

$$p^*(z, t) = P \exp(s t) \exp(2\pi i \frac{z}{\lambda})$$

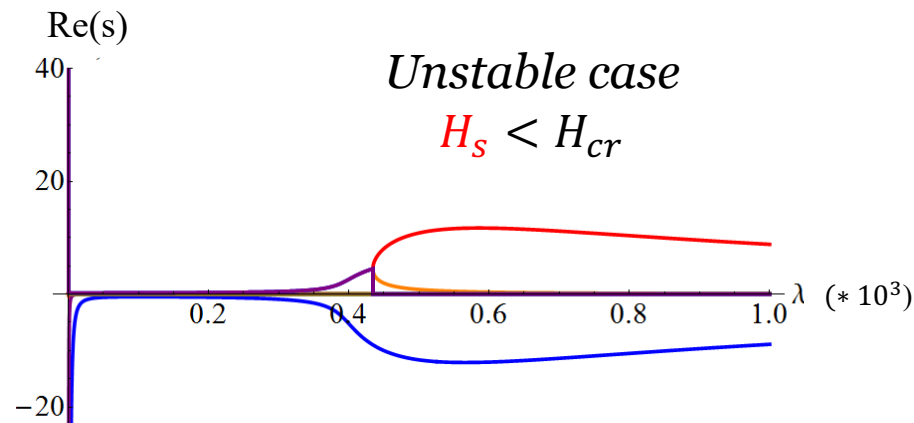
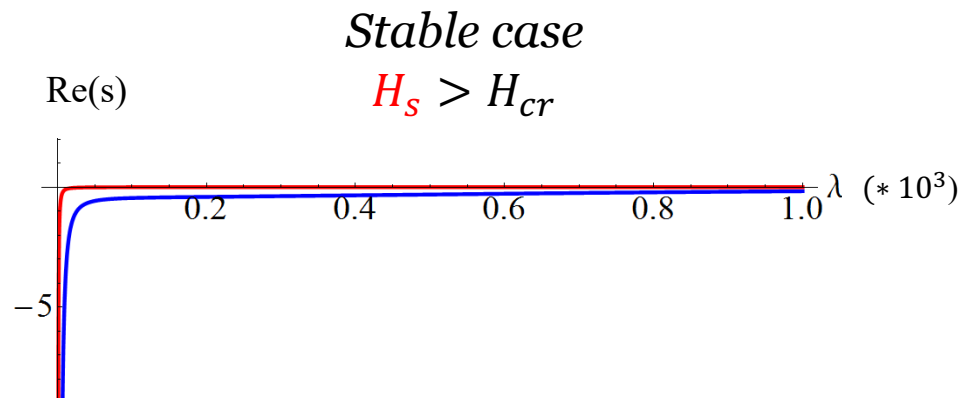
- Onset of strain localization: sign of $\text{Re}(s)$.

Unstable $s > 0$



Bifurcation analysis

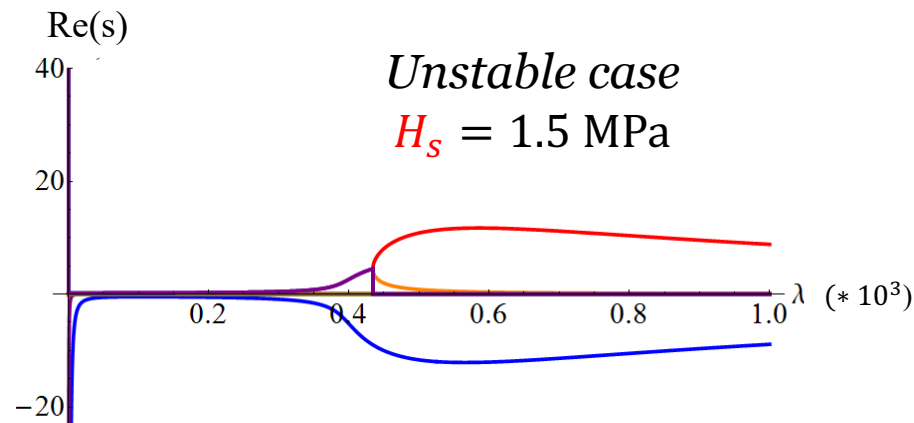
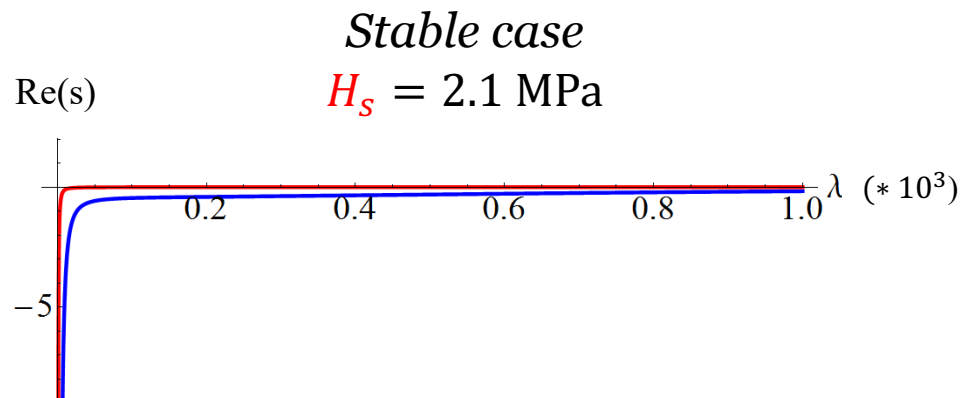
Critical value of the hardening H_{cr} for the onset of localization.



Example of bifurcation analysis

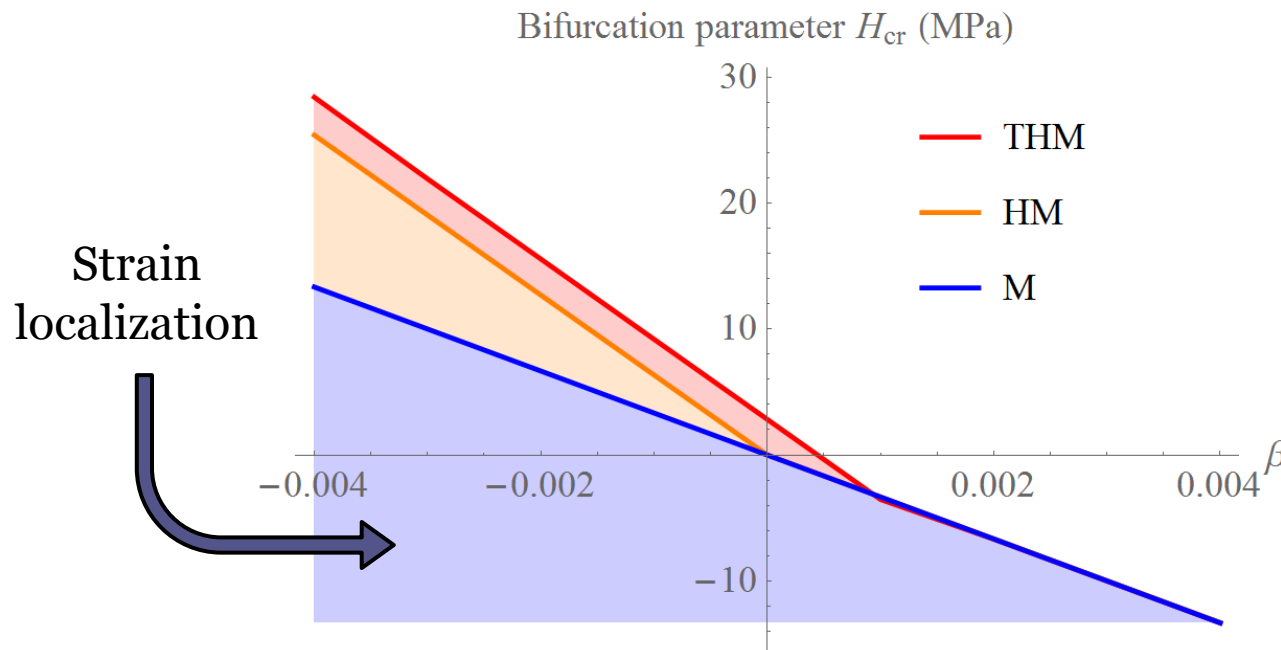
Constitutive parameters refer to a fault zone at 7 km depth.

⇒ Bifurcation for a critical hardening modulus of $H_{cr} = 2 \text{ MPa} > 0$



Onset of strain localization

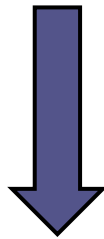
- Dilatancy β strongly influences the bifurcation.



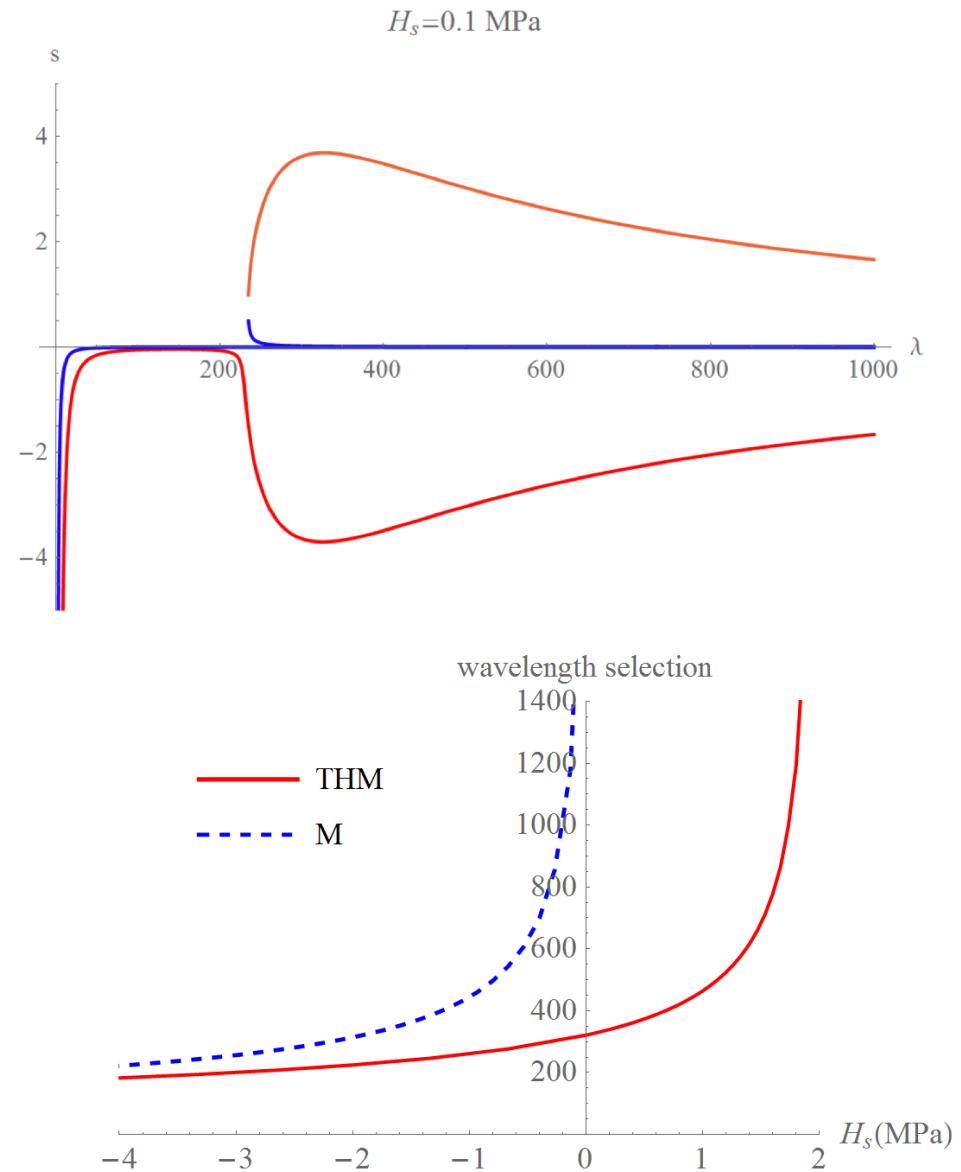
- HM couplings destabilize the system for contractant materials ($\beta < 0$).
- THM couplings make the system unstable for dilatant materials in the hardening regime.

Wavelength selection

Evolution of the maximum
with the hardening modulus

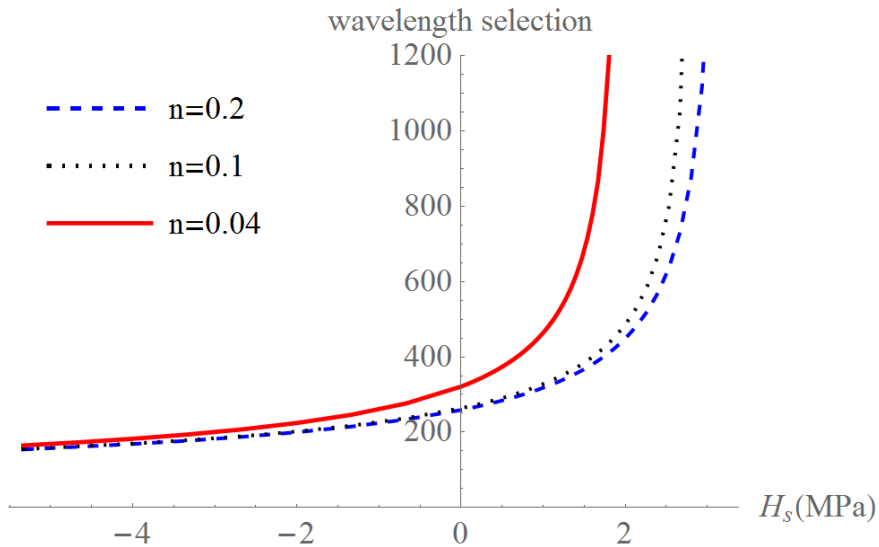


Evolution of the shear band thickness
in function of the hardening modulus

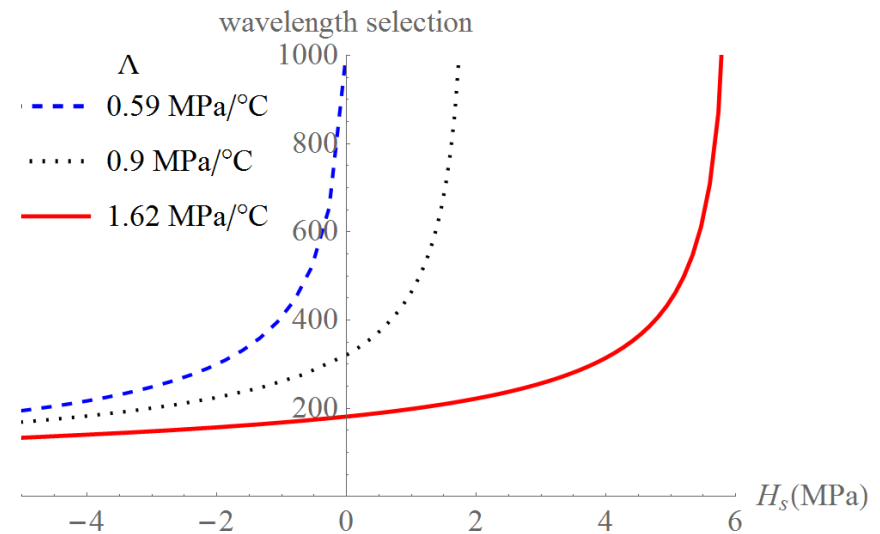


Dominant parameters on strain localization ?

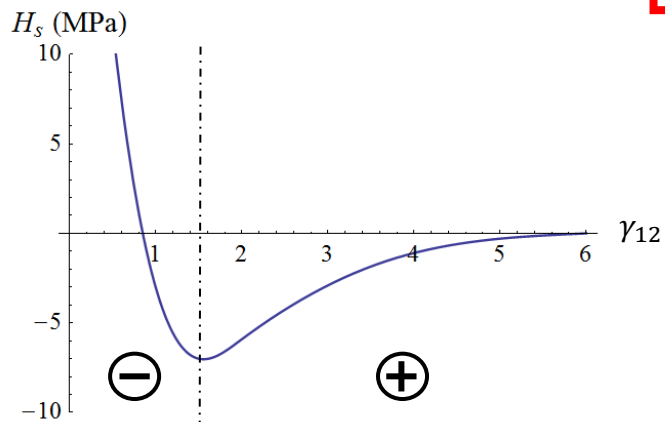
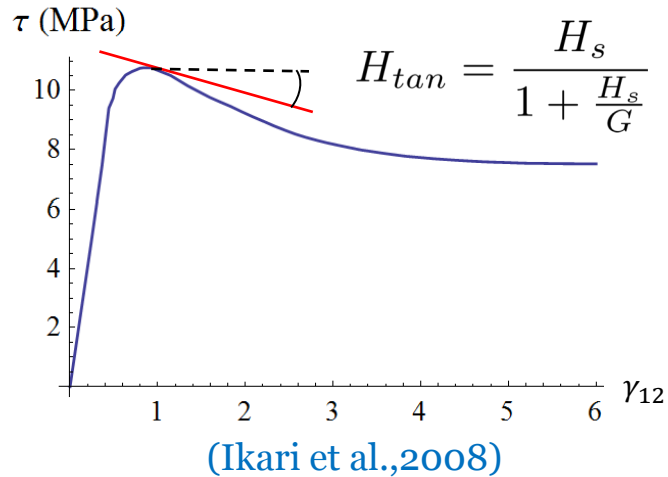
Porosity



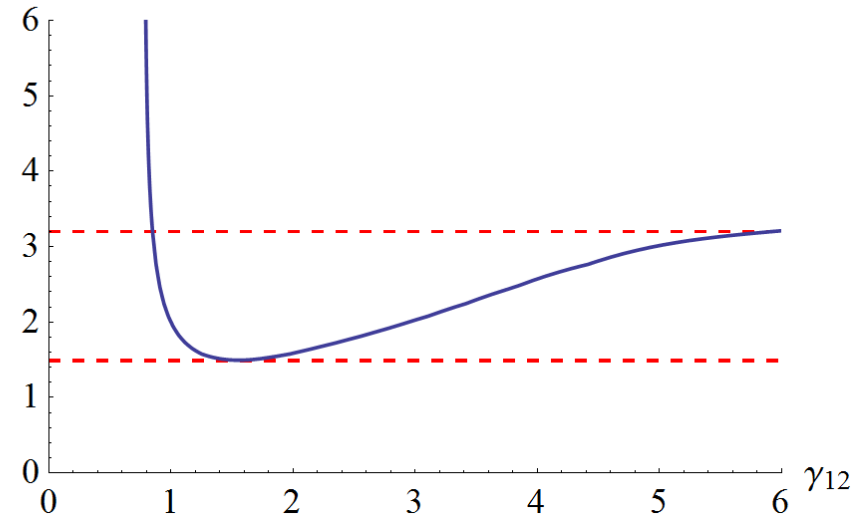
Thermal pressurization coefficient Λ



Shear band thickness evolution

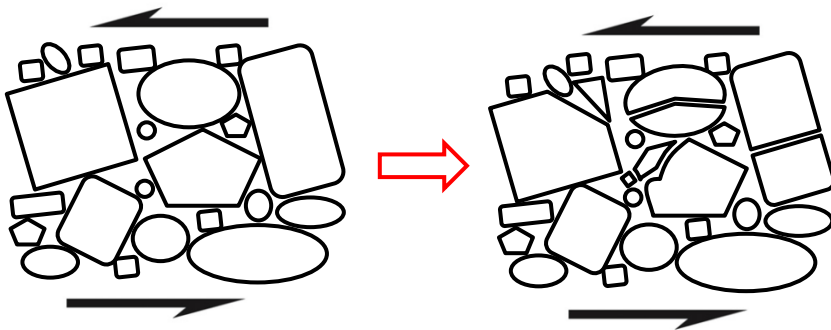


shear band width (mm)

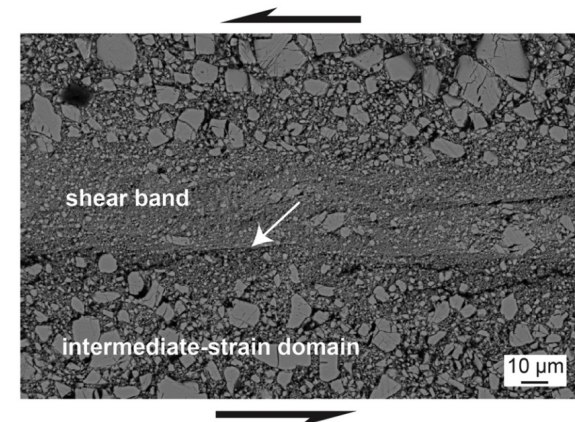


Grain cataclasis

- Grain crushing at high mean stresses.
- Smaller grains inside the shear band.
- 30 % reduction after a strain of 15 (Gu et al., 1994).



Shear experiments on Dolomite
(Smith et al., 2015)



Effect of grain size reduction

- Exponential evolution of D_{50} with the total shear strain γ_{12} :

$$D(\gamma_{12}) = (D_0 - D_{fin}) \exp\left(-\frac{\gamma_{12}}{\gamma_c}\right) + D_{fin}$$

- Progressive decrease** of the shear band thickness.

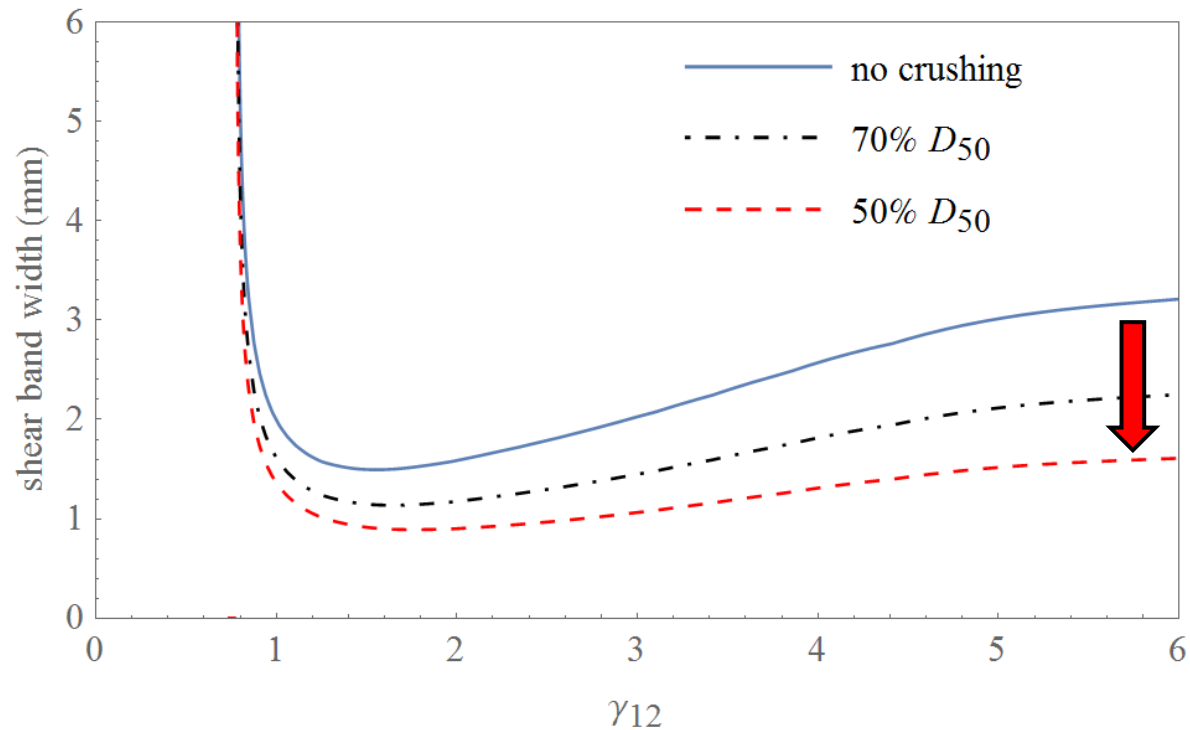


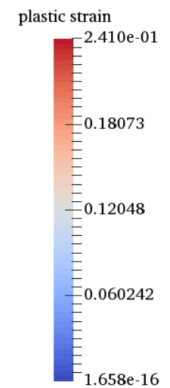
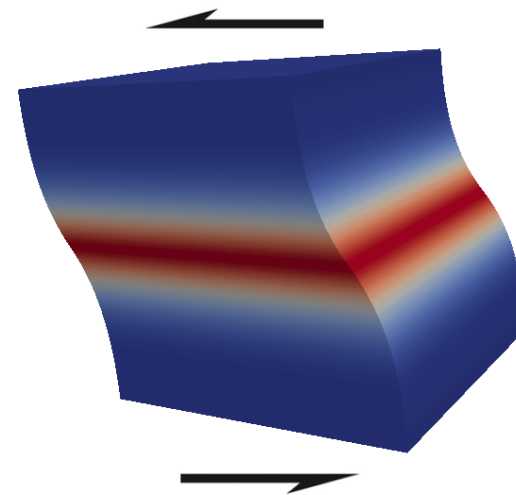
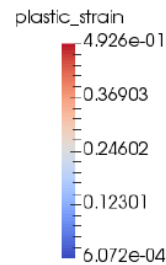
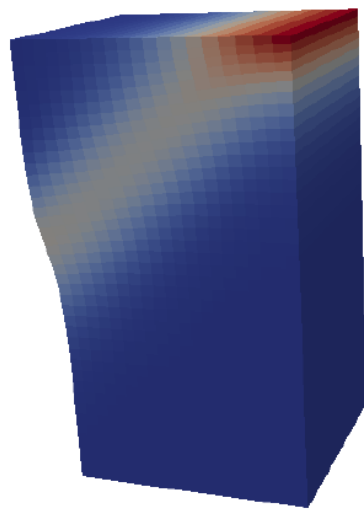
Table of contents

- ❑ Fault mechanics and strain localization
- ❑ Mathematical modeling
- ❑ Linear Stability Analysis
- ❑ **Finite Element simulations**
- ❑ Conclusions and perspectives

FEM analysis of a Cosserat THM model

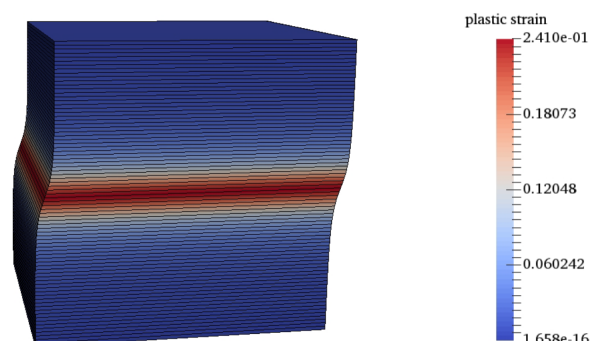
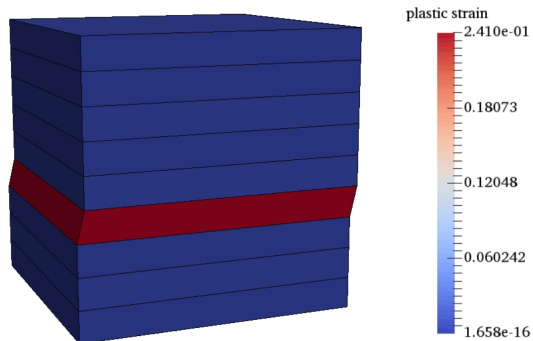
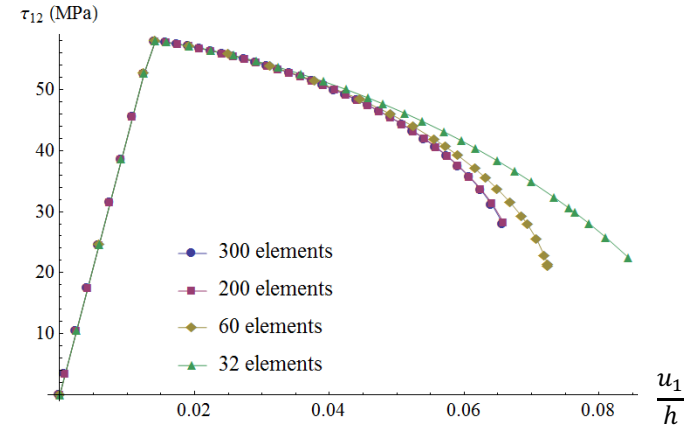
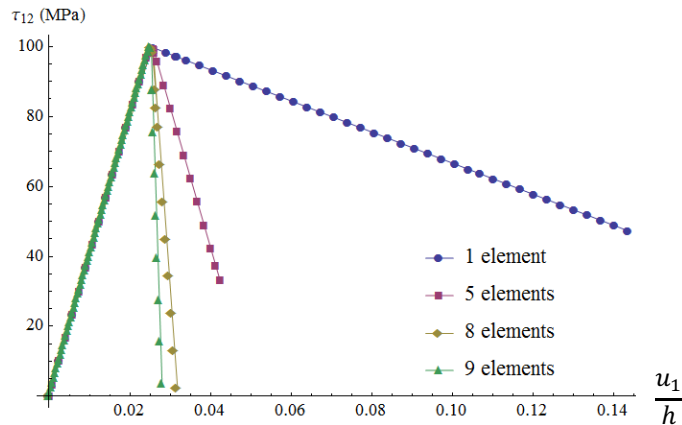


Implementation of a 3D fully coupled elasto-plastic Cosserat continuum:
Redback/Moose.



Mesh independency

- Finite shear band thickness.
- Correct representation of **weakening** and **energy dissipation**.



Finite element implementation



- Weak form of the momentum balance equations:

$$-\int_{\Omega} \tau_{ij} \psi_{i,j} d\Omega + \int_{\partial\Omega_{\Sigma}} \tau_{ij} n_j \psi_i dS = 0$$

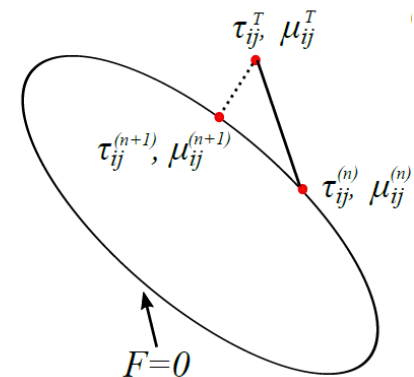
$$-\int_{\Omega} \mu_{ij} \psi_{i,j} d\Omega + \int_{\partial\Omega_{\Sigma}} \mu_{ij} n_j \psi_i dS - \int_{\Omega} \varepsilon_{ijk} \tau_{jk} \psi_i d\Omega = 0$$

- Weak form of energy and fluid mass balance equations:

$$\int_{\Omega} \dot{p} \psi d\Omega + c_{hy} \left(\int_{\Omega} p_{,i} \psi_{,i} d\Omega - \int_{\partial\Omega} p_{,i} n_i \psi dS \right) - \Lambda \int_{\Omega} \dot{T} \psi d\Omega + \frac{1}{\beta^*} \int_{\Omega} \varepsilon_v \psi d\Omega = 0$$

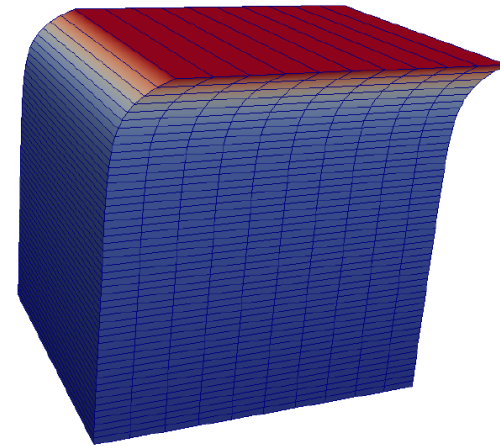
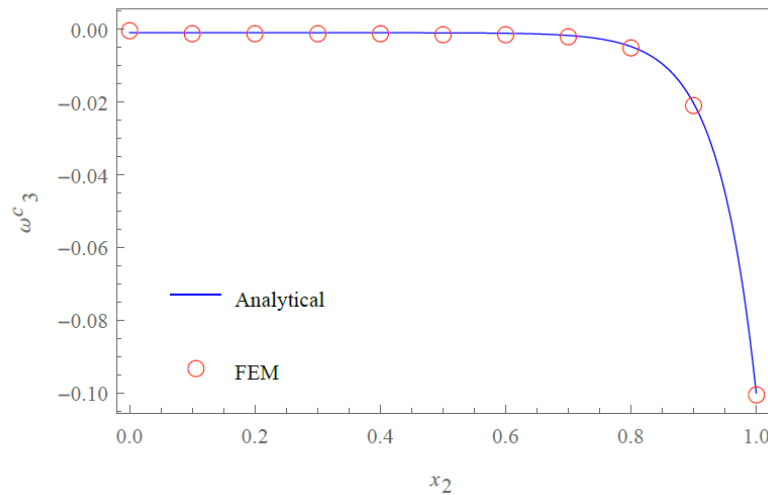
$$\int_{\Omega} \dot{T} \psi d\Omega + c_{th} \left(\int_{\Omega} T_{,i} \psi_{,i} d\Omega - \int_{\partial\Omega} T_{,i} n_i \psi dS \right) - \frac{1}{\rho C} \int_{\Omega} (\tau_{ij} \dot{\gamma}_{ij}^p + \mu_{ij} \dot{\kappa}_{ij}^p) \psi_{,i} d\Omega = 0$$

- ψ and ψ_i are linear Lagrange test functions.
- Incremental plastic constitutive law is integrated using a return map algorithm (Godio et al., 2016).

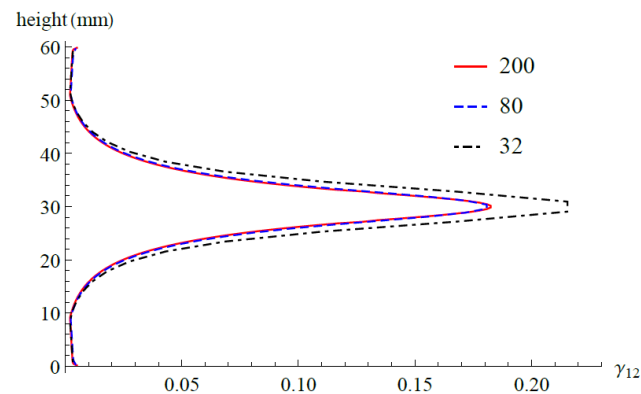


Validation

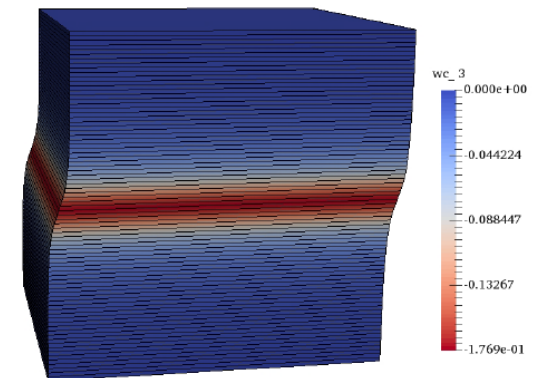
- Elastic test: Boundary layer effect (Vardoulakis & Sulem, 1995).



- Plastic test: Shearing of a layer with von Mises yield surface (Godio et al., 2016).

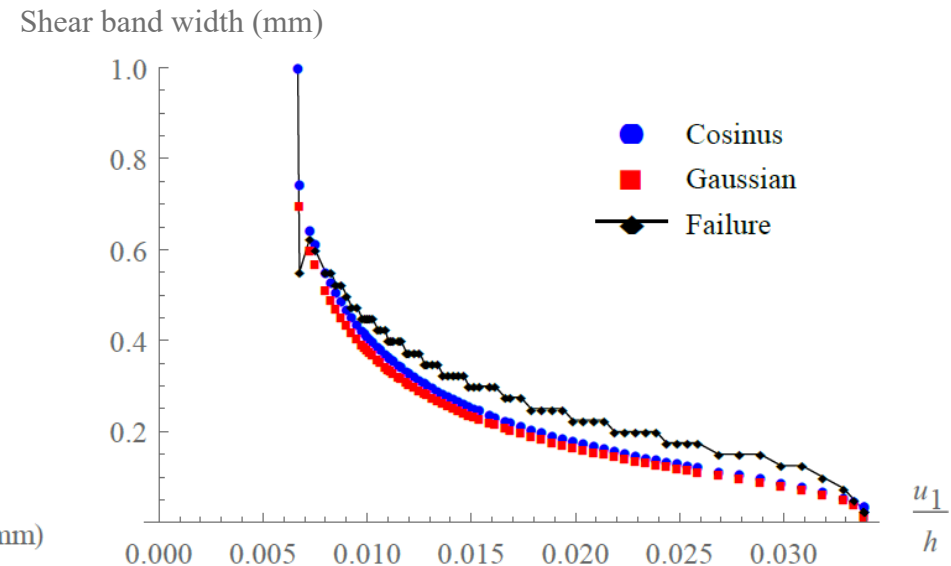
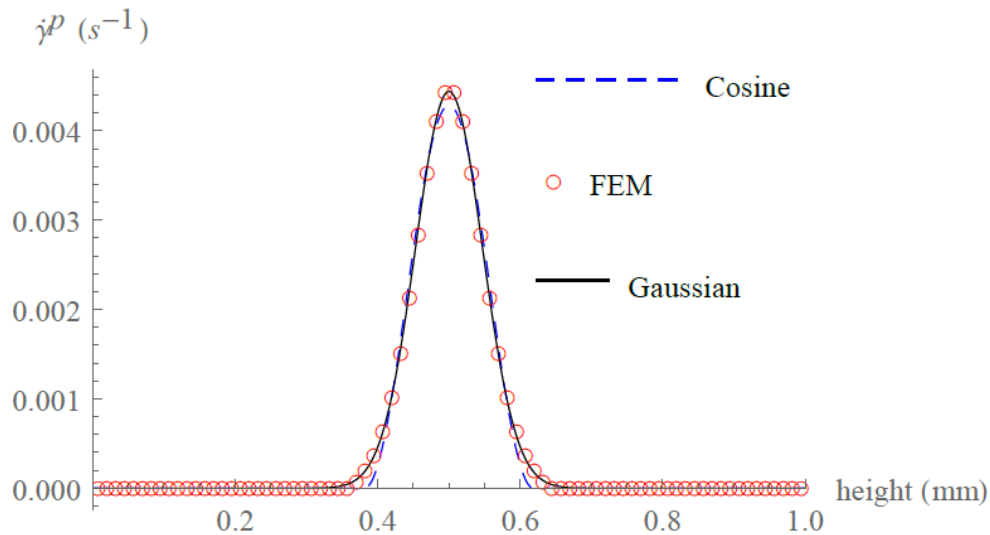


- Various couplings...



How to evaluate the shear band thickness ?

➔ Fitting of the **plastic strain rate** profile.



Linear softening

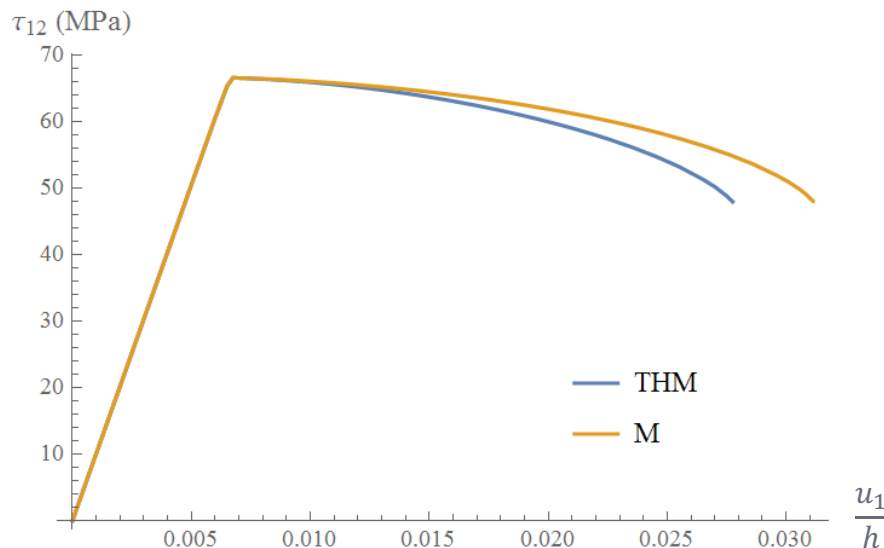
- Drucker-Prager yield surface with linear evolution of the friction coefficient.

$$F = \tau + \mu \sigma' - c$$

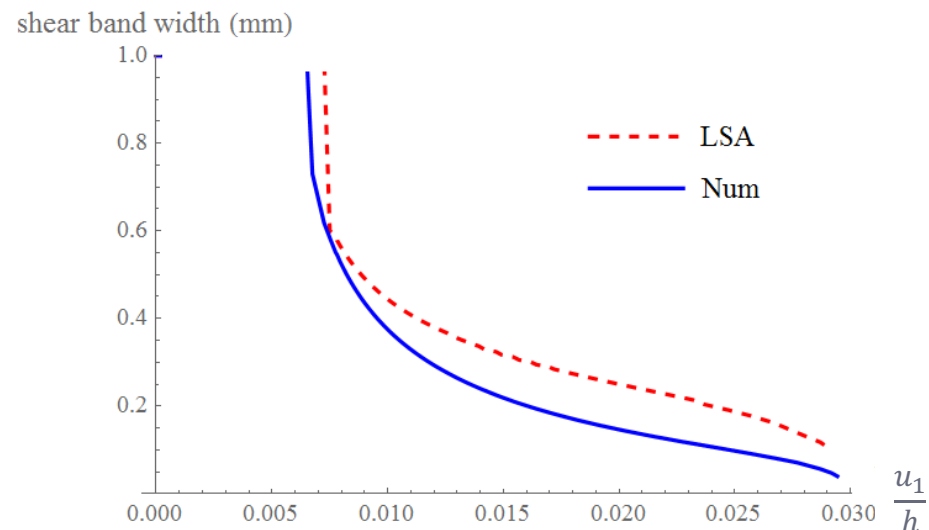
$$\mu = \mu_0(1 + H_s q)$$

$$\dot{q} = \dot{\gamma}^p$$

Load-displacement



Localisation for THM



thickness obtained by LSA compatible with FEM.

Exponential softening

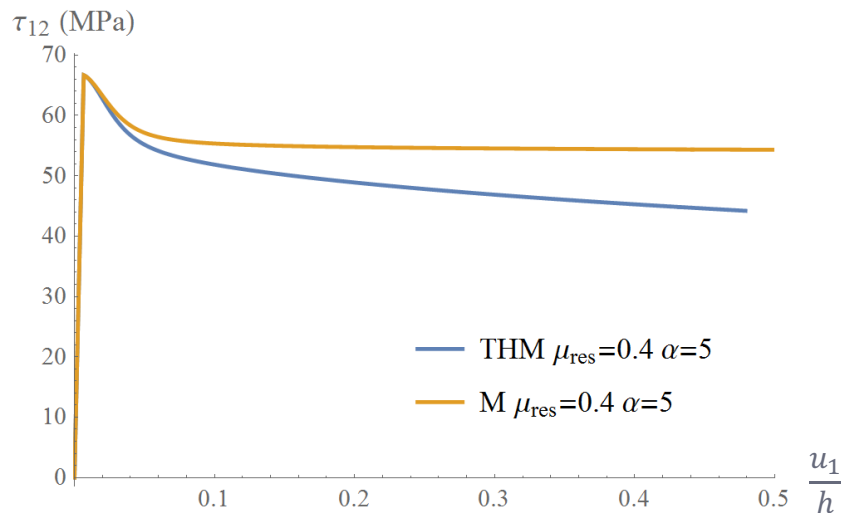
- Drucker-Prager yield with exponential decrease of the friction coefficient.

$$F = \tau + \mu \sigma' - c$$

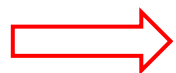
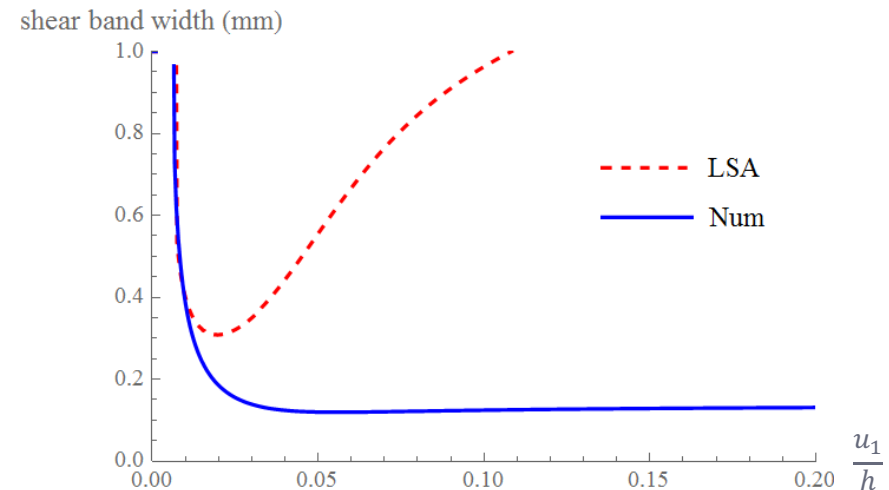
$$\mu = \mu_0 \left(1 + \frac{\Delta\mu}{\mu_0} \exp(H_s q) \right)$$

$$\dot{q} = \dot{\gamma}^p$$

Load-displacement



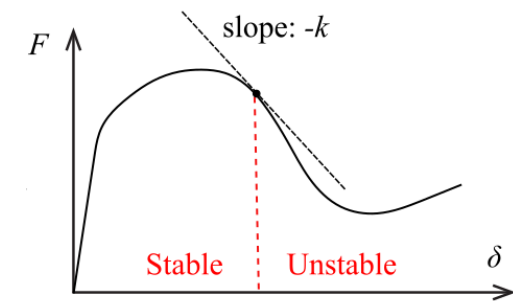
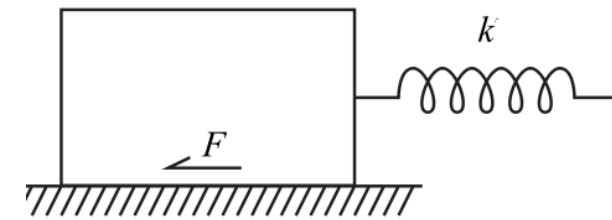
Evolution of the friction coefficient



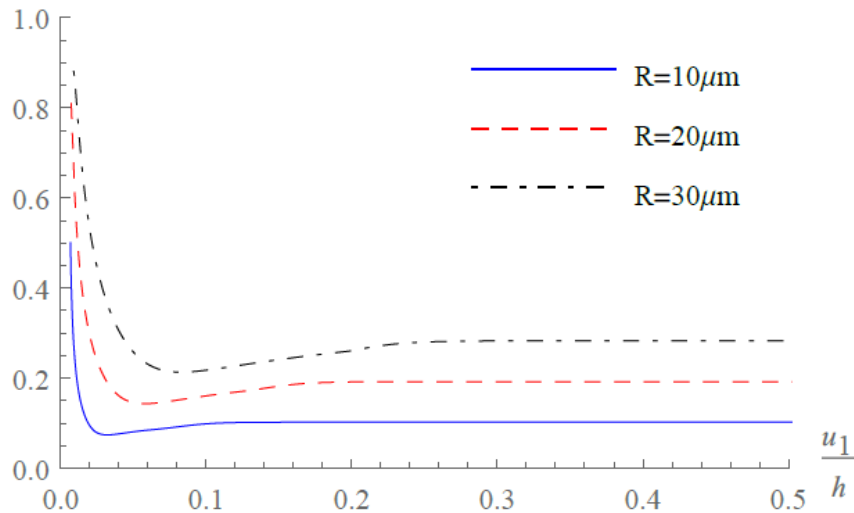
Limit of the LSA: Perturbed fields dominate after first stages of deformation.

Effect of the microstructure

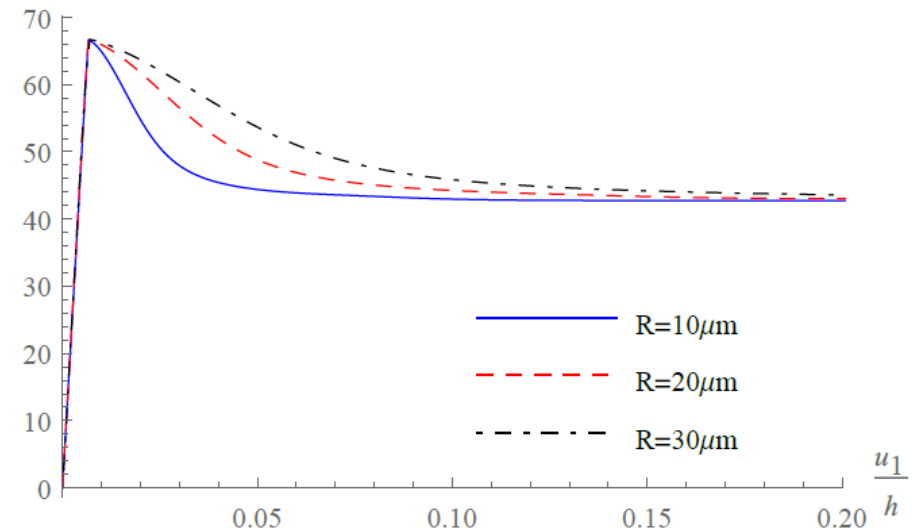
- The grain size affects the shear band thickness.
- And thus, the stress-strain diagram.



shear band width (mm)

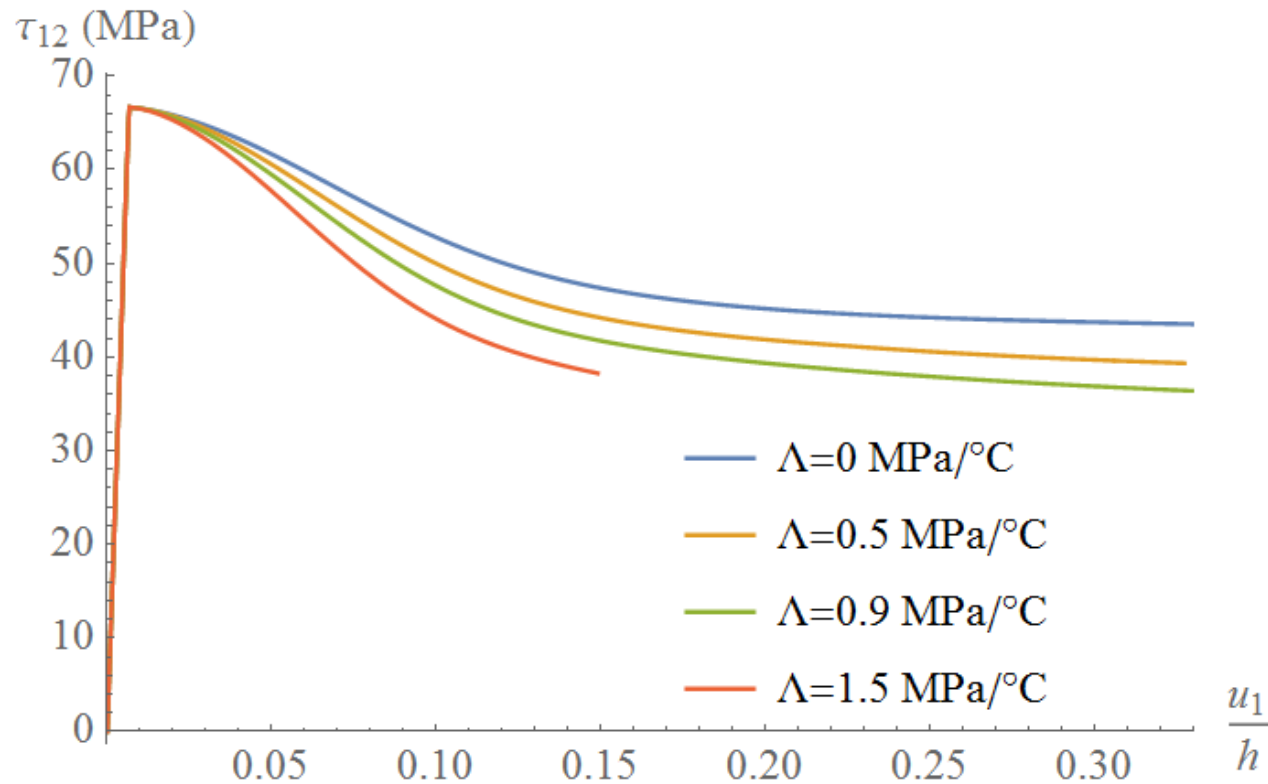


τ_{12} (MPa)



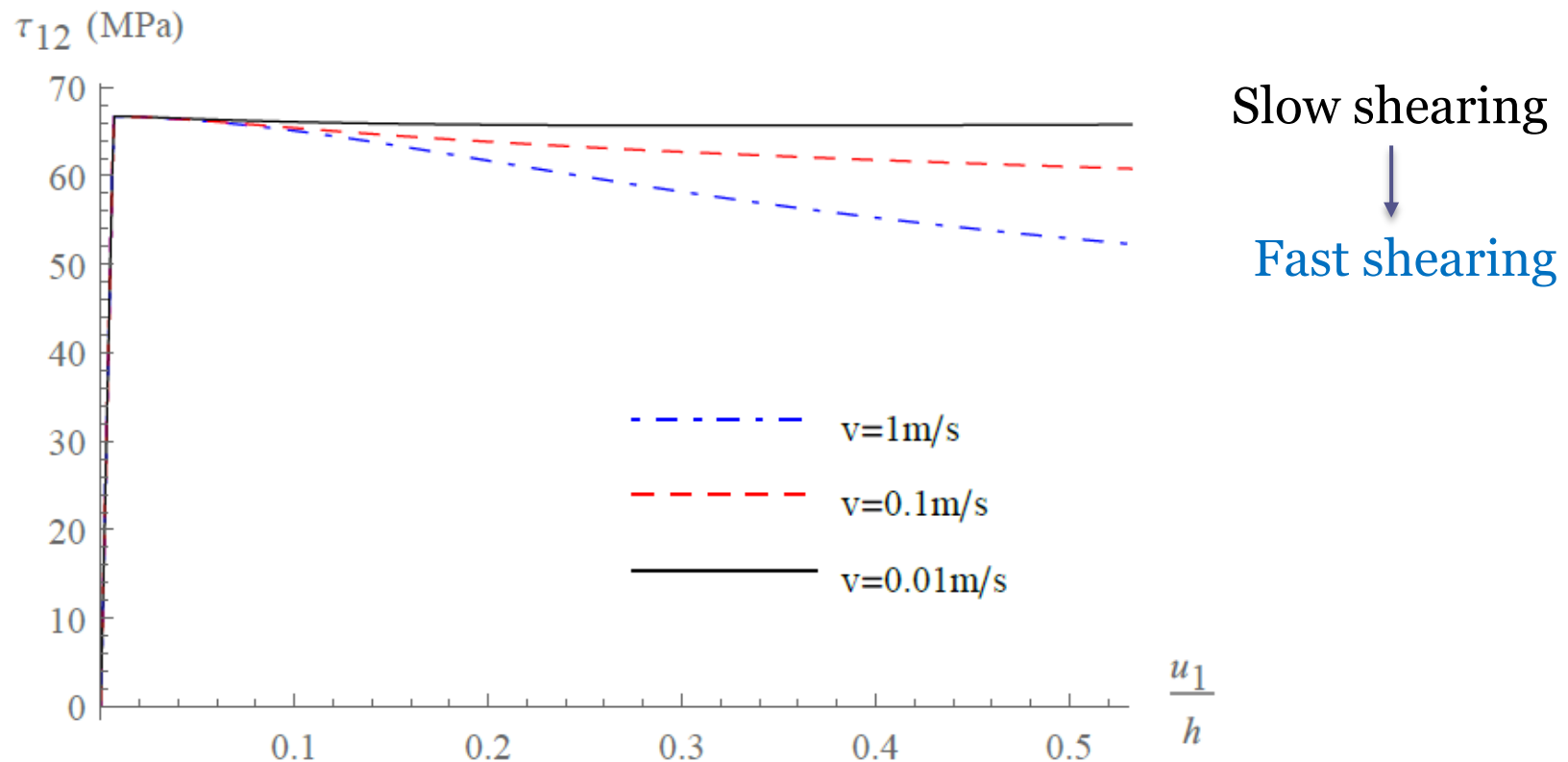
Effect of the thermal pressurization

A higher thermal pressurization coefficient Λ induces a stronger softening (and a narrower slip zone).



Apparent rate-dependency

Despite the use of a rate-independent constitutive law (perfect plasticity here).

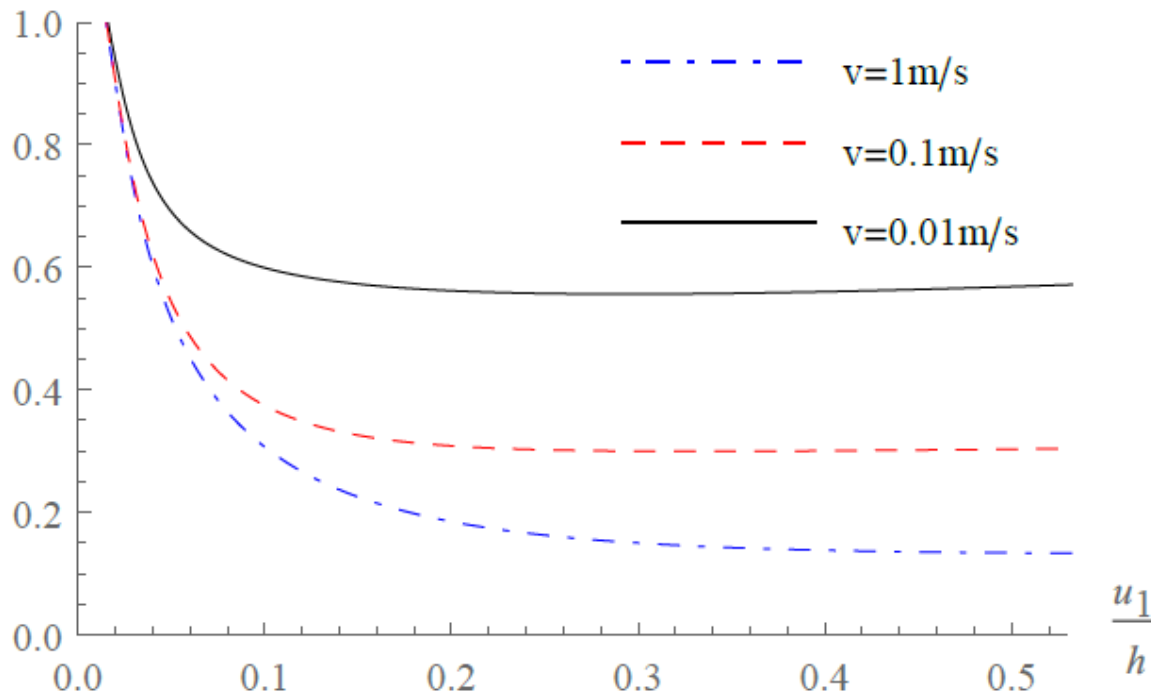


Induced by the THM couplings.

Rate-dependency of strain localization

Diffusion processes change the localization thickness

shear band width (mm)



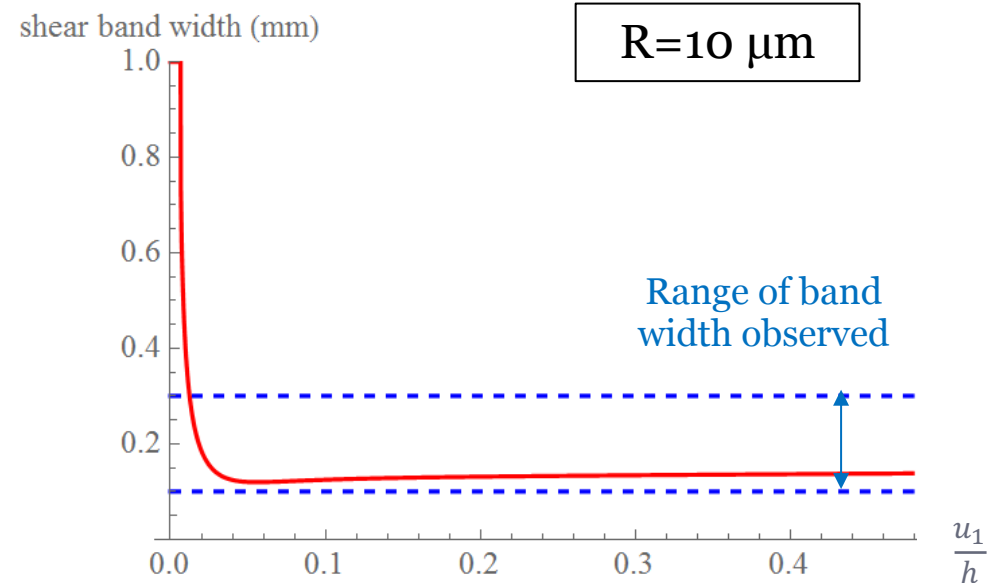
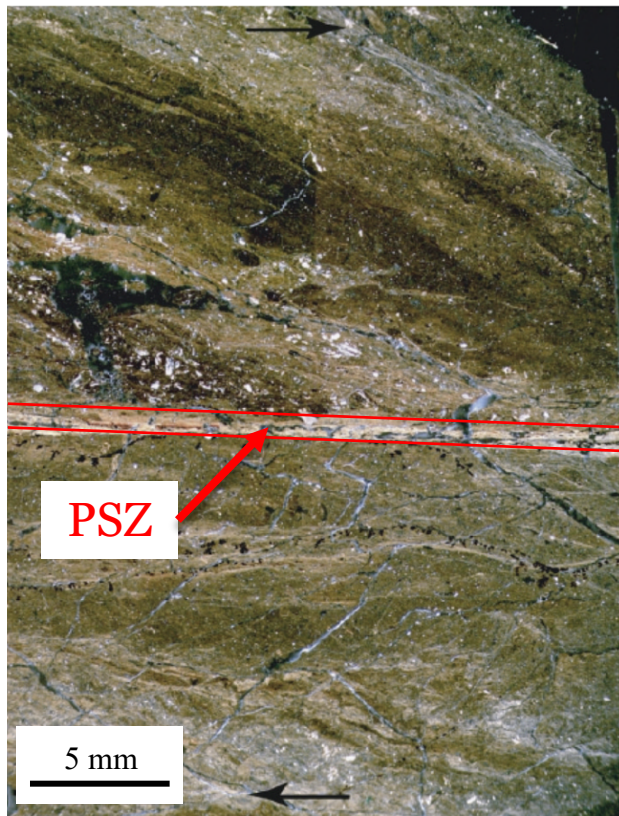
Slow shearing



Fast shearing

Comparison with field observations

Example of the Punchbowl fault: Principal Slip Zone (PSZ) of **100-300 μm** thick.



Shear band thickness prediction consistent with field observations

Conclusions (1/2)

- The use of the **Cosserat** continuum is appropriate for modeling fault gouge **microstructure** and obtaining **finite thickness of the localized zone**. Thus, **Energy dissipation** is correctly quantified.
- The onset of strain localization and the evolution of the shear band thickness are investigated with LSA.
- Development of a robust numerical tool for fully coupled Cosserat THM systems.
- Post Bifurcation regime and shear band thickness evolution are investigated with FEM considering THM couplings.

Conclusions (2/2)

- The thickness obtained by LSA are comparable to FEM as long as the assumptions for linearization around a homogeneous state are valid.
- **Rate dependency** appears (in this model) at the macro-scale as a result of an evolving shear band thickness due to thermal pressurization and diffusion processes.
- The predicted thickness of the localized fault zone is **in agreement with field observations**.

Perspectives

Short term

- Extension of the analysis to account for other weakening processes active in fault zones during seismic slip (grain cataclasis, chemical effects, flash heating, ...).
- Assessment from the numerical analyses of the fracture energy and compare it with experimental and field observations.
- Comparison or merge with a rate dependent plastic model.

Long term

- Need for advanced laboratory experiments and field observations to constrain the model.
- Large scale modelling.

Thank you for your attention !